

Divergent Box Integral 4: $I_4^{D=4-2\epsilon}(0, 0, p_3^2, p_4^2; s_{12}, s_{23}; 0, 0, 0, 0)$

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The result for the box (see **figure**) in the unphysical region ($s_{12} < 0, s_{23} < 0, p_2^2 < 0, p_4^2 < 0$) is[1]

$$\begin{aligned} I_4^{D=4-2\epsilon}(0, 0, p_3^2, p_4^2; s_{12}, s_{23}; 0, 0, 0, 0) &= \frac{\mu^{2\epsilon}}{s_{12}s_{23}} \\ &\times \left[\frac{2}{\epsilon^2} \left((-s_{12})^{-\epsilon} + (-s_{23})^{-\epsilon} - (-p_3^2)^{-\epsilon} - (-p_4^2)^{-\epsilon} \right) + \frac{1}{\epsilon^2} \left((-p_3^2)^{-\epsilon} (-p_4^2)^{-\epsilon} \right) / (-s_{12})^{-\epsilon} \right. \\ &\left. - 2 \operatorname{Li}_2 \left(1 - \frac{p_3^2}{s_{23}} \right) - 2 \operatorname{Li}_2 \left(1 - \frac{p_4^2}{s_{23}} \right) - \ln^2 \left(\frac{-s_{12}}{-s_{23}} \right) \right] + \mathcal{O}(\epsilon). \end{aligned}$$

In ref. [1] the auxiliary box function for two adjacent external off-shell lines, is introduced, related to the six dimensional box,

$$\begin{aligned} \operatorname{LS}_{-1}^{2mh}(s_{12}, s_{23}; p_3^2, p_4^2) &= -\operatorname{Li}_2 \left(1 - \frac{p_3^2}{s_{23}} \right) - \operatorname{Li}_2 \left(1 - \frac{p_4^2}{s_{23}} \right) - \frac{1}{2} \ln^2 \left(\frac{-s_{12}}{-s_{23}} \right) + \frac{1}{2} \ln \left(\frac{-s_{12}}{-p_3^2} \right) \ln \left(\frac{-s_{12}}{-p_4^2} \right) \\ &\quad + \left[\frac{1}{2} (s_{12} - p_3^2 - p_4^2) + \frac{p_3^2 p_4^2}{s_{23}} \right] I_3^{3m}(s_{12}, p_3^2, p_4^2), \end{aligned} \quad (1)$$

where I_3^{3m} is the scalar triangle integral with three lines off shell. This integral vanishes in the appropriate ‘back-to-back’ kinematic limit. They also introduce a version of this box function with I_3^{3m} removed,

$$\tilde{\operatorname{LS}}_{-1}^{2mh}(s_{12}, s_{23}; p_3^2, p_4^2) = -\operatorname{Li}_2 \left(1 - \frac{p_3^2}{s_{23}} \right) - \operatorname{Li}_2 \left(1 - \frac{p_4^2}{s_{23}} \right) - \frac{1}{2} \ln^2 \left(\frac{-s_{12}}{-s_{23}} \right) + \frac{1}{2} \ln \left(\frac{-s_{12}}{-p_3^2} \right) \ln \left(\frac{-s_{12}}{-p_4^2} \right). \quad (2)$$

Thus we may write

$$\begin{aligned} I_4^{D=4-2\epsilon}(0, 0, p_3^2, p_4^2; s_{12}, s_{23}; 0, 0, 0, 0) &= \frac{\mu^{2\epsilon}}{s_{12}s_{23}} \\ &\times \left[\frac{2}{\epsilon^2} \left((-s_{12})^{-\epsilon} + (-s_{23})^{-\epsilon} - (-p_3^2)^{-\epsilon} - (-p_4^2)^{-\epsilon} \right) + \frac{1}{\epsilon^2} \left((-p_3^2)^{-\epsilon} (-p_4^2)^{-\epsilon} \right) / (-s_{12})^{-\epsilon} \right. \\ &\left. - \ln \left(\frac{-s_{12}}{-p_3^2} \right) \ln \left(\frac{-s_{12}}{-p_4^2} \right) + 2 \tilde{\operatorname{LS}}_{-1}^{2mh}(s_{12}, s_{23}; p_3^2, p_4^2) \right] + \mathcal{O}(\epsilon). \end{aligned}$$

The result for the six dimensional box is

$$I_4^{D=6}(0, 0, p_3^2, p_4^2; s_{12}, s_{23}; 0, 0, 0, 0) = \frac{s_{23} \text{LS}_{-1}^{2mh}(s_{12}, s_{23}; p_3^2, p_4^2)}{p_3^2 p_4^2 - s_{23} s_{13}}.$$

See the file on [notation](#).

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References

- [1] Z. Bern, L. J. Dixon and D. A. Kosower, Nucl. Phys. B **412**, 751 (1994) [[arXiv:hep-ph/9306240](#)]
- [2] G. Duplancic and B. Nizic, Eur. Phys. J. C **20**, 357 (2001) [[arXiv:hep-ph/0006249](#)]